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Approximation on the rotation group SO(3)

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Abstract

In this paper we study the approximation on rotation group SO(3), we consider the partial sum, Féjer and Jackson-type operators and obtain the approximation theorems in $L_p(1 \le p \le +\infty)$ respectively. ©2017 All rights reserved.

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1. Introduction

Many results of approximation are based on Euclid spaces or their compact subsets. Periodic approximation is based on compact group $\{\exp(ix)\}$, whereas matrix group U(n) is the generalization of $\{\exp(ix)\}$. We know homomorphism between SU(2) and rotation group SO(3), which has many applications in Physics and Chemistry. Some approximation problems on compact groups have been studied since in 1920s Peter and Weyl [7] proved the approximation theorem on compact group, that is, the irreducible character generates a dense subspace of the space of continuous classes function. For instance, Gong [2] studied the basic problems of Fourier analysis on unitary and rotation groups, including the degree of convergence of Abel sum based on Poisson kernel. Zheng et al. (see [11, 12]) studied the polynomial approximation on compact Lie groups. Cartwright et al. [1] studied Jackson's theorem for compact connected Lie groups is different from on the classical cases, Eulid spaces. For example, Riemann-Lebesgue lemma is not necessarily true in $L_p(G)(1 \le p \le 4/3)$, with G being a compact Lie group (see [8]). In this paper, we study the approximation of the partial sums, Féjer and Jackson-type operators on SO(3) in $L_p(1 \le p \le +\infty)$.

Let $G = SO(3) = \{x \in GL(3, \mathbb{R}), x^T x = E, \det x = 1\}$ be the rotation group, where $GL(n, \mathbb{R})$ is the group of invertible real $(n \times n)$ -matrices. For $1 \le p < +\infty$, $L_p(G) = \{f : ||f||_p = [\int_G |f(x)|^p d\mu(x)]^{1/p} < +\infty\}$, and

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 μ is the normalized Harr measure on G, conveniently, we write $d\mu(x)$ as dx in this paper. When using Euler angles, the Harr integral of a function f on SO(3) reads as

$$\int_{\mathrm{SO}(3)} f(\mathbf{x}) d\mu(\mathbf{x}) = \frac{1}{8\pi^2} \int_0^{2\pi} d\alpha \int_0^{\pi} \sin\beta d\beta \int_0^{2\pi} d\gamma f(\alpha, \beta, \gamma)$$

An element $\mathbf{x} \in SO(3)$ is identified with a point in the projective spaces S_{π} of the closed ball in \mathbb{R}^3 of radius π by $\mathbf{x} \to \omega \mathbf{r}$, satisfying $\mathbf{xr} = \mathbf{r}$, and $||\mathbf{r}|| = 1$, where \mathbf{r} is the rotation axis and $\omega \in [0, \pi]$ is the rotation angle of \mathbf{x} . (ϕ, θ, φ) denotes the direction angles of rotation axis \mathbf{r} and ω the rotation angle, then there hold following formulas (see [9]) between (\mathbf{r}, ω) and Euler angles (α, β, γ),

$$\cos\frac{\omega}{2} = \cos\frac{\beta}{2}\cos\frac{\alpha+\gamma}{2}, \ \varphi = -\frac{\sin\frac{\beta}{2}\sin\frac{\alpha-\gamma}{2}}{\sin\frac{\omega}{2}}, \ \theta = \frac{\sin\frac{\beta}{2}\cos\frac{\alpha-\gamma}{2}}{\sin\frac{\omega}{2}}, \ \varphi = \frac{\cos\frac{\beta}{2}\sin\frac{\alpha+\gamma}{2}}{\sin\frac{\omega}{2}}.$$
 (1.1)

We know a function f on SO(3) only depends on the rotation angle of the argument coincide exactly with the class function-functions that are conjugacy classes, i.e., $f(\mathbf{x}) = f(\mathbf{y}\mathbf{x}\mathbf{y}^{-1})$ for all $\mathbf{x}, \mathbf{y} \in SO(3)$. In other words, for any function f on SO(3) there is a uniquely determined $\tilde{f} : [0, \pi] \to \mathbb{C}$ such that $f(\mathbf{x}) = \tilde{f}(\omega(\mathbf{x}))$. The Haar integral for such a class function reads as

$$\int_{\mathrm{SO}(3)} \mathbf{f}(\mathbf{x}) d\mu(\mathbf{x}) = \int_{\mathrm{SO}(3)} \tilde{\mathbf{f}}(\omega(\mathbf{x})) d\mu(\mathbf{x}) = \frac{2}{\pi} \int_0^{\pi} \tilde{\mathbf{f}}(\omega) \sin^2 \frac{\omega}{2} d\omega.$$
(1.2)

We make the convention that if f is a class function on SO(3) and if no confusion occurs, we subsequently drop the tilde. Denote by

$$S_{N}f(\alpha,\beta,\gamma) = \sum_{l=0}^{N} \sum_{m,n=-l}^{l} \sqrt{2l+1}C_{m,n}^{l}D_{m,n}^{l}(\alpha,\beta,\gamma),$$

the partial sum operators, where $(\alpha, \beta, \gamma) \in [0, 2\pi] \times [0, \pi] \times [0, 2\pi]$ are Euler angles and

$$C_{m,n}^{l} = \frac{\sqrt{2l+1}}{8\pi^2} \int_0^{2\pi} d\alpha \int_0^{\pi} \sin\beta d\beta \int_0^{2\pi} d\gamma f(\alpha,\beta,\gamma) \overline{D_{m,n}^{l}(\alpha,\beta,\gamma)},$$

and for $\mathbf{x} \in SO(3)$,

$$D_{\mathfrak{m},\mathfrak{n}}^{\mathfrak{l}}(\mathbf{x}) = D_{\mathfrak{m},\mathfrak{n}}^{\mathfrak{l}}(\alpha,\beta,\gamma) = e^{-i\mathfrak{m}\alpha}P_{\mathfrak{m},\mathfrak{n}}^{\mathfrak{l}}(\cos\beta)e^{-i\mathfrak{n}\gamma}, -\mathfrak{l}\leqslant\mathfrak{m},\mathfrak{n}\leqslant\mathfrak{l},$$

the function $P_{m,n}^{l}$ is given by

$$P_{m,n}^{l}(u) = C(1-u)^{\frac{n-m}{2}}(1+u)^{-(\frac{n+m}{2})}\frac{d^{l-m}}{du^{l-m}}[(1-u)^{l-n}(1+u)^{l+n}],$$
with $C = \frac{(-1)^{l-n}i^{n-m}}{2^{l}}\sqrt{\frac{(l+m)!}{(l-n)!(l+n)!(l-m)!}}.$

2. Theorems and their proofs

We start with the partial sum operator,

$$\begin{split} S_{N}(f, \mathbf{x}) &= S_{N}f(\alpha, \beta, \gamma) \\ &= \sum_{l=0}^{N} \sum_{m,n=-l}^{l} \sqrt{2l+1} C_{m,n}^{l} D_{m,n}^{l}(\alpha, \beta, \gamma) \end{split}$$

$$\begin{split} &= \sum_{l=0}^{N} \sum_{m,n=-l}^{l} \frac{2l+1}{8\pi^2} \int_{0}^{2\pi} d\zeta \int_{0}^{\pi} \sin\eta d\eta \int_{0}^{2\pi} d\xi f(\zeta,\eta,\xi) \overline{D_{m,n}^{l}(\zeta,\eta,\xi)} D_{m,n}^{l}(\alpha,\beta,\gamma) \\ &= \sum_{l=0}^{N} \frac{2l+1}{8\pi^2} \int_{0}^{2\pi} d\zeta \int_{0}^{\pi} \sin\eta d\eta \int_{0}^{2\pi} d\xi f(\zeta,\eta,\xi) \sum_{m,n=-l}^{l} \overline{D_{m,n}^{l}(\zeta,\eta,\xi)} D_{m,n}^{l}(\alpha,\beta,\gamma) \\ &= \sum_{l=0}^{N} (2l+1) \int_{SO(3)} f(\mathbf{g}) Tr(D^{l}(\mathbf{g}^{-1}\mathbf{x})) d\mathbf{g} \\ &= \sum_{l=0}^{N} (2l+1) f * \chi_{l}(\mathbf{x}) = f * D_{N}(\mathbf{x}), \end{split}$$

where $D^{l}(\mathbf{x}) = \text{matrix}(D^{l}_{m,n})$ is the (2l+1)-dimensional irreducible representation of G, $\chi_{l}(\mathbf{x}) = \text{Tr}[D^{l}(\mathbf{x})]$ is the character of irreducible representation, $D_N(\mathbf{x}) = \sum_{l=0}^{N} (2l+1)\chi_l(\mathbf{x})$ is Dirichlet kernel, and the convolution of two functions f and g is defined by

$$(\mathbf{f} * \mathbf{g})(\mathbf{x}) = \int_{\mathbf{G}} \mathbf{f}(\mathbf{g}) \mathbf{g}(\mathbf{g}^{-1}\mathbf{x}) d\mathbf{g}.$$

Next orthogonal properties of D^1 and χ_1 are well-known (see [3]),

$$\frac{2j_1+1}{8\pi^2} \int_0^{2\pi} d\alpha \int_0^{\pi} \sin\beta d\beta \int_0^{2\pi} d\gamma \overline{D_{m_1,n_1}^{j_1}(\alpha,\beta,\gamma)} D_{m_2,n_2}^{j_2}(\alpha,\beta,\gamma) = \delta_{j_1,j_2} \delta_{m_1,m_2} \delta_{n_1,n_2},$$
$$\frac{1}{\pi} \int_0^{\pi} \overline{\chi_{j_1}(\omega)} \chi_{j_2}(\omega) (1-\cos\omega) d\omega = \delta_{j_1,j_2}.$$

Furthermore, we discuss the properties of Dirichlet kernel D_N . Noting that $\chi_1(\mathbf{x}) = Tr[D^1(\mathbf{x})] = \frac{\sin(2l+1)\frac{\omega}{2}}{\sin\frac{\omega}{2}}$ with ω as (1.1) being the rotation angle of \mathbf{x} and using Euler formula $e^{i\theta} = \cos \theta + i \sin \theta$, by calculation, we can get the following.

Lemma 2.1.

$$D_{N}(\mathbf{x}) = \frac{\sin(N+1)\omega \cdot \cos\frac{\omega}{2}}{2\sin^{3}\frac{\omega}{2}} - \frac{(N+1)\cdot\cos(N+1)\omega}{\sin^{2}\frac{\omega}{2}},$$
$$\|D_{N}\|_{L_{1}(SO(3))} = O(N) + O(\ln N), \quad \|S_{N}\| = O(N).$$

For $f \in L_p(G)$, $1 \leq p \leq +\infty$, we define the best approximation degree by

$$\mathsf{E}_{\mathsf{N}}(\mathsf{f})_{\mathfrak{p}} = \sup_{\mathsf{T}_{\mathsf{N}}\in\Pi_{\mathsf{N}}} \|\mathsf{f}-\mathsf{T}_{\mathsf{N}}\|_{\mathfrak{p}},$$

where Π_N is the set of trigonometric polynomial of degree at most N, which needs not to be an integer, and when $p = +\infty$, $L_{\infty}(G)$ is replaced by continuous functions space C(G) and the p-norm is replaced by sup-norm, we will not repeat this later.

Thanks to the accuracy of S_N for the trigonometric polynomial, that is $S_N(T_n) = T_n(n \leq N)$, $T_n \in \Pi_N$, subsequently we have the following.

Theorem 2.2. *If* $f \in L_p(SO(3)), 1 \leq p \leq \infty$ *, then*

$$\|\mathbf{f} - \mathbf{S}_{N}(\mathbf{f})\|_{p} \leq CN \cdot \mathbf{E}_{N}(\mathbf{f})_{p}$$

with C > 0 being an absolute constant.

Further, we consider the Féjer kernel

$$F_{N}(\mathbf{x}) = \frac{D_{0}(\mathbf{x}) + D_{1}(\mathbf{x}) + \dots + D_{N}(\mathbf{x})}{N+1}$$

Using Lemma 2.1, we know

$$F_{N}(\mathbf{x}) = \frac{1}{N+1} \left[\frac{\cos \frac{\omega}{2}}{2\sin^{3} \frac{\omega}{2}} \sum_{k=0}^{N} \sin(k+1)\omega - \frac{1}{\sin^{2} \frac{\omega}{2}} \sum_{k=0}^{N} (k+1)\cos(k+1)\omega \right].$$

Note that

$$\sum_{k=0}^{N} \sin(k+1)\omega = \operatorname{Im}\left[\sum_{k=0}^{N} (e^{i\omega})^{k+1}\right] = \left(\sin\frac{N+2}{2}\omega\right) \cdot \frac{\sin(\frac{N+1}{2}\omega)}{\sin\frac{\omega}{2}},$$

and

$$\begin{split} &\sum_{k=0}^{N} (k+1)\cos(k+1)\omega \\ &= \mathrm{Im}[\sum_{k=0}^{N} (e^{\mathrm{i}\omega})^{k+1}]' \\ &= \frac{N+2}{2}(\cos\frac{N+2}{2}\omega) \cdot \frac{\sin\frac{N+1}{2}\omega}{\sin\frac{\omega}{2}} + \frac{N+1}{2}(\cos\frac{N+1}{2}\omega) \cdot \frac{\sin\frac{N+2}{2}\omega}{\sin\frac{\omega}{2}} - \frac{1}{2}\cos\frac{\omega}{2} \cdot \frac{\sin\frac{N+1}{2}\omega \cdot \sin\frac{N+2}{2}\omega}{\sin^2\frac{\omega}{2}}, \end{split}$$

we have

$$\begin{split} \mathsf{F}_{\mathsf{N}}(\mathbf{x}) &= \frac{1}{2(\mathsf{N}+1)} \frac{\cos\frac{\omega}{2}}{\sin^{2}\frac{\omega}{2}} \cdot \frac{\sin\frac{\mathsf{N}+1}{2}\omega \cdot \sin\frac{\mathsf{N}+2}{2}\omega}{\sin^{2}\frac{\omega}{2}} - \frac{1}{\sin^{2}\frac{\omega}{2}} \Big[\frac{\mathsf{N}+2}{2(\mathsf{N}+1)} (\cos\frac{\mathsf{N}+2}{2}\omega) \cdot \frac{\sin\frac{\mathsf{N}+1}{2}\omega}{\sin\frac{\omega}{2}} \\ &+ \frac{1}{2} (\cos\frac{\mathsf{N}+1}{2}\omega) \cdot \frac{\sin\frac{\mathsf{N}+2}{2}\omega}{\sin\frac{\omega}{2}} - \frac{1}{2(\mathsf{N}+1)} \cos\frac{\omega}{2} \cdot \frac{\sin\frac{\mathsf{N}+1}{2}\omega \cdot \sin\frac{\mathsf{N}+2}{2}\omega}{\sin^{2}\frac{\omega}{2}} \Big] \\ &:= \mathsf{I}_{1}(\omega) + \mathsf{I}_{2}(\omega) + \mathsf{I}_{3}(\mathsf{t}) + \mathsf{I}_{4}(\omega). \end{split}$$
(2.1)

Lemma 2.3.

$$\|F_N\|_{L_1(SO(3))} = O(\ln N).$$

From $D_N(\mathbf{x}) = \sum_{l=0}^{N} (2l+1)\chi_l(\mathbf{x})$, and the orthogonal properties of D^1 and χ_l , we have

$$\int_{\mathrm{SO}(3)} \mathsf{F}_{\mathsf{N}}(\mathbf{x}) d\mathbf{x} = 1.$$

We consider the next Féjer convolution operator

$$\sigma_{\mathsf{N}}(\mathbf{f},\mathbf{x}) = \int_{\mathsf{SO}(3)} \mathsf{F}_{\mathsf{N}}(\mathbf{y}) \mathsf{f}(\mathbf{y}^{-1}\mathbf{x}) d\mathbf{y}.$$

Clearly,

$$|f(\mathbf{x}) - \sigma_{\mathsf{N}}(f, \mathbf{x})| \leq \int_{\mathsf{SO}(3)} |f(\mathbf{x}) - f(\mathbf{y}^{-1}\mathbf{x})| |\mathsf{F}_{\mathsf{N}}(\mathbf{y})| d\mathbf{y} \leq \int_{\mathsf{SO}(3)} \omega(f, |\mathbf{y}|) |\mathsf{F}_{\mathsf{N}}(\mathbf{y})| d\mathbf{y}, \tag{2.2}$$

where |y| = dist(y, e) and e is the identity element of SO(3). For $x, y \in SO(3)$, $\text{dist}(x, y) = \omega(y^{-1}x) = \omega(x^{-1}y)$ is defined as the rotational angle of the rotation $y^{-1}x$ (or its inverse $x^{-1}y$).

Remark 2.4. The rotational angle of a rotation $\mathbf{x} \in SO(3)$ is also defined as $|\mathbf{x}| = \arccos \frac{1}{2}[\operatorname{Tr}(\mathbf{x}) - 1]$, in terms of Euler angles $|\mathbf{x}| = |\mathbf{x}(\alpha, \beta, \gamma)| = 2 \arccos \left(\cos \frac{\beta}{2} \cos \frac{\alpha + \gamma}{2}\right) = \omega(\mathbf{x}) = \omega$ (see [4]).

Remark 2.5. $\omega(f,t)_p = \sup\{||f(x) - f(y^{-1}x)||_p : d(x,y) \leq t\}(1 \leq p \leq \infty)$, is the first-order moduli of smoothness. Higher order moduli $\omega_r(f,t)$ (see Definition 2.8) will also be employed later. We need only the inequality $\omega_r(f,\lambda t)_p \leq (1+\lambda)^r \omega_r(f,t)_p (r \in \mathbb{N})$, some other properties can be found in [1, 5].

From Remarks 2.4 and 2.5 and (1.2) we have

$$\begin{split} |f(\mathbf{x}) - \sigma_{N}(f, \mathbf{x})| &\leqslant \omega(f, \frac{1}{N}) \cdot \int_{SO(3)} (1 + N|\mathbf{y}|) |F_{N}(\mathbf{y})| d\mathbf{y} \\ &\leqslant \omega(f, \frac{1}{N}) \cdot \frac{2}{\pi} \int_{0}^{\pi} (1 + N\omega) |F_{N}(\omega)| \sin^{2} \frac{\omega}{2} dt \\ &= \omega(f, \frac{1}{N}) \cdot \frac{2}{\pi} \int_{0}^{\pi} (1 + N\omega) \left[|I_{1}(\omega)| + |I_{2}(\omega)| + |I_{3}(\omega)| + |I_{4}(\omega)| \right] \sin^{2} \frac{\omega}{2} dt. \end{split}$$

By (2.1) and subdividing the integration interval $[0, \pi] = [0, \frac{\pi}{N+2}] \cup [\frac{\pi}{N+2}, \pi]$, we get

$$\int_{0}^{\pi} (1 + N\omega) |I_{j}(\omega)| \sin^{2} \frac{\omega}{2} d\omega = O(\ln(N)), j = 1, 2, 3, 4$$

Thus, by using Lemma 2.3 and (2.2), we have the following result.

Theorem 2.6. *If* $f \in L_p(SO(3))$, $1 \le p \le +\infty$, *then*,

$$\|f - \sigma_N(f)\|_p = O(\ln(N) \cdot \omega(f, \frac{1}{N}))_p, \ 1 \leqslant p \leqslant \infty.$$

Generally, we conclude the result as follows.

Theorem 2.7. If $f \in L_p(SO(3))$, $1 \le p \le +\infty$, and $V_N \in \Pi_N$ satisfies

(i): $V_N * T_N = T_N, \forall T_N \in \Pi_N$, (ii): $\|V_N\|_{L_1(G)} \leq K$, with K being an absolute constant, then for $L_N(f) = V_N * f$, $\|f - L_N(f)\|_p \leq (1+K)E_N(f)_p$.

Proof. The proof is trivial. The kernel V_N , which satisfies the conditions of Theorem 2.7, can be found in [6].

Next we give the Jackson-type approximation theorem.

We first give the definition of the r-th moduli of smoothness of function f on the rotation group. As classical case, for $\mathbf{x}, \mathbf{h} \in G = SO(3)$, let

$$(\Delta_h^r f)(x) = \sum_{j=0}^r (-1)^{r-j} \binom{r}{j} f(h^{-j}x)$$

be the r-th difference of f at the point **x**.

Definition 2.8. Let $f \in L_p(G)(1 \le p \le \infty)$, for any integer $r \ge 1$ and t > 0, write

$$\omega_{\mathbf{r}}(\mathbf{f},\mathbf{t})_{\mathbf{p}} = \sup \left\{ \|\Delta_{\mathbf{h}}^{\mathbf{r}}\mathbf{f}\|_{\mathbf{p}} : \mathbf{h} \in \mathbf{G}, \mathbf{d}(\mathbf{e},\mathbf{h}) = |\mathbf{h}| \leq \mathbf{t} \right\}.$$

We can also write

$$\omega_{r}(f,t)_{p} = \sup \left\{ \|\Delta_{expH}^{r}f\|_{p} : H \in \mathbf{g}, \|H\| \leq t \right\}$$

where g is the Lie algebras of G, and $\|\cdot\|$ is the norm induced by Killing inner product \ll on g (see [10]).

Let T be a maximal torus of SO(3) given by (see [7])

$$\mathsf{T} = \left\{ \left(\begin{array}{ccc} \cos \omega & \sin \omega & 0 \\ -\sin \omega & \cos \omega & 0 \\ 0 & 0 & 1 \end{array} \right) \middle| \omega \in \mathbb{R}/2\pi\mathbb{Z} \right\},\$$

and t denotes its respective Lie algebras also called Cartan subalgebras, in fact given by

$$\mathfrak{t} = \left\{ \left(\begin{array}{ccc} 0 & \omega & 0 \\ -\omega & 0 & 0 \\ 0 & 0 & 0 \end{array} \right) \middle| \omega \in \mathbb{R}/2\pi\mathbb{Z} \right\}.$$

For $\mathbf{u} \in SO(3)$, let $\phi_{N,s}(\mathbf{u}) = \lambda_{N,s}^{-1} [\frac{\sin(N+\frac{1}{2})\omega}{\sin\frac{1}{2}\omega}]^{2s} = \phi_{N,s}(\omega)$, satisfying $\int_{SO(3)} \phi_{N,s}(\mathbf{u}) d\mathbf{u} = 1$, where $\omega = \omega(\mathbf{u})$ is the rotation angle of \mathbf{u} , and $s \ge 2 + r/2$. As the classical case, we have $\lambda_{N,s} = O(N^{2s-3})$, and

$$\int_{SO(3)} |\mathbf{u}|^k \phi_{N,s}(\mathbf{u}) d\mathbf{u} = \frac{2}{\pi} \int_0^{\pi} \phi_{N,s}(\omega) \sin^2 \frac{\omega}{2} d\omega = O(N^{-k}), k = 0, 1, \cdots, 2s - 4.$$
(2.3)

Following the thought of classical case, we construct the Jackson-type approximate operator as follows, which is different from [1].

$$J_{N}(f, \mathbf{x}) = \int_{G} \phi_{N,s}(\mathbf{u}) \sum_{j=1}^{r} (-1)^{j-1} \binom{r}{j} f(\mathbf{u}^{-j}\mathbf{x}) d\mathbf{u}$$

By the Weyl integral formula (see [10]),

$$\begin{split} J_{N}(f,\mathbf{x}) &= \int_{G} \varphi_{N,s}(\mathbf{u}) \sum_{j=1}^{r} (-1)^{j-1} {r \choose j} f(\mathbf{u}^{-j}\mathbf{x}) d\mathbf{u} \\ &= \sum_{j=1}^{r} (-1)^{j-1} {r \choose j} \int_{G} \varphi_{N,s}(\mathbf{u}) f(\mathbf{u}^{-j}\mathbf{x}) d\mathbf{u} \\ &= |W(G)|^{-1} \sum_{j=1}^{r} (-1)^{j-1} {r \choose j} \int_{T} d(\mathbf{v}) \varphi_{N,s}(\mathbf{v}) \int_{G/T} f(\mathbf{g}\mathbf{v}^{-j}\mathbf{g}^{-1}\mathbf{x}) dT d(G/T) \\ &= |W(G)|^{-1} |Q|^{-1} \sum_{j=1}^{r} (-1)^{j-1} {r \choose j} \int_{Q} |D(\mathbf{H})|^{2} \varphi_{N,s}(\mathbf{H}) \int_{G/T} f(\mathbf{g}\exp(-j\mathbf{H})\mathbf{g}^{-1}\mathbf{x}) dgT dH \\ &= |W(G)|^{-1} |Q|^{-1} \int_{Q} \int_{G/T} \sum_{j=1}^{r} (-1)^{j-1} {r \choose j} \frac{1}{j} |D(\mathbf{H}/j)|^{2} \varphi_{N,s}(\mathbf{H}/j) f(\mathbf{g}\exp(-\mathbf{H})\mathbf{g}^{-1}\mathbf{x}) dgT dH \\ &= |W(G)|^{-1} |Q|^{-1} \int_{Q} \int_{G/T} \sum_{j=1}^{r} (-1)^{j-1} {r \choose j} \frac{1}{j} |D(\mathbf{H}/j)|^{2} \varphi_{N,s}(\mathbf{H}/j) f(\mathbf{g}\exp(-\mathbf{H})\mathbf{g}^{-1}\mathbf{x}) dgT dH \\ &= |W(G)|^{-1} |Q|^{-1} \int_{Q} \int_{G/T} \sum_{j=1}^{r} (-1)^{j-1} {r \choose j} \frac{1}{j} |D(\mathbf{H}/j)|^{2} \varphi_{N,s}(\mathbf{H}/j) f(\mathbf{g}\exp(-\mathbf{H})\mathbf{g}^{-1}\mathbf{x}) dgT dH \\ &= |W(G)|^{-1} |Q|^{-1} \int_{Q} \int_{G/T} |D(\mathbf{H})|^{2} K_{N}(\exp \mathbf{H}) f(\mathbf{g}\exp(-\mathbf{H})\mathbf{g}^{-1}\mathbf{x}) dgT dH \\ &= |W(G)|^{-1} \int_{T} \int_{G/T} d(\mathbf{v}) K_{N}(\mathbf{v}) f(\mathbf{g}\mathbf{v}^{-1}\mathbf{g}^{-1}\mathbf{x}) dgT d\mathbf{v} \\ &= \int_{G} K_{N}(\mathbf{u}) f(\mathbf{u}^{-1}\mathbf{x}) d\mathbf{u}, \end{split}$$

where

$$K_{N}(\mathbf{u}) = K_{N}(\mathbf{v}) = K_{N}(\exp \mathbf{H}) = \sum_{j=1}^{r} (-1)^{j-1} {r \choose j} \frac{1}{j} \phi_{N,s}(\mathbf{H}/j) \left| \frac{D(\mathbf{H}/j)}{D(\mathbf{H})} \right|^{2}$$

$$= \sum_{j=1}^{r} (-1)^{j-1} {r \choose j} \frac{1}{j} \phi_{N,s}(\omega/j) \left| \frac{\sin(\omega/2j)}{\sin \omega/2} \right|^{2},$$
(2.4)

with $\mathbf{u} = \mathbf{g}\mathbf{v}\mathbf{g}^{-1}$, $\mathbf{v} \in T$, and $\mathbf{v} = \exp \mathbf{H}$, $\mathbf{H} \in \mathfrak{t}$, and $D(\mathbf{H}) = 2\mathfrak{i}\sin\frac{\omega}{2}$, $D(\mathbf{H}/\mathfrak{j}) = 2\mathfrak{i}\sin(\omega/2\mathfrak{j})$, $d(\mathbf{v}) = 4\sin^2\frac{\omega}{2} := d(\omega)$, $\exp \mathfrak{i}\mathfrak{g} \to G$ is the exponential map, $\exp |\mathfrak{t} : \mathfrak{t} \to T$.

Weyl group W = W(G) is defined by W = N/T, with $N = \{g \in G | gTg^{-1} = T\}$, T as above is maximal torus of G, |W(G)| = |W(SO(3))| = 2 denotes the order of W. Q is called fundamental domain centered at the origin, which satisfies $\exp(Q) = T$, here in fact $Q = \mathbf{t} \cong [-\pi, \pi]$, $|Q| = 2\pi$ denotes the volume of Q.

Lemma 2.9. If $f \in L_p(SO(3))$, $1 \le p \le +\infty$, then $J_N(f) \in \Pi_N$.

Proof. In fact

$$J_{N}(f, \mathbf{x}) = (K_{N} * f)(\mathbf{x}) = \int_{G} K_{N}(\mathbf{u}) f(\mathbf{u}^{-1}\mathbf{x}) d\mathbf{u} = \int_{G} K_{N}(\mathbf{x}\mathbf{u}^{-1}) f(\mathbf{u}) d\mathbf{u}.$$

Clearly, from (2.4), $K_N(\mathbf{u})$ is a trigonometric polynomial of degree N, which has not to be an integer.

We write $K_N(\mathbf{u}) = \sum_{l=0}^{N} \sum_{m,n=-l}^{l} c_{m,n}^l D_{m,n}^l(\mathbf{u})$, and note that (see [4]) $D_{m,n}^l(\mathbf{u}_1\mathbf{u}_2) = \sum_{k=-l}^{l} D_{m,k}^l(\mathbf{u}_1) D_{k,n}^l(\mathbf{u}_2)$, $D_{m,n}^l(\mathbf{u}^{-1}) = [D_{m,n}^l(\mathbf{u})]^{-1} = [\overline{D}_{m,n}^l(\mathbf{u})] = \overline{D}_{n,m}^l(\mathbf{u})$, $J_N(\mathbf{f}, \mathbf{x}) = \int_{G} K_N(\mathbf{x}\mathbf{u}^{-1}) \mathbf{f}(\mathbf{u}) d\mathbf{u} = \sum_{l=0}^{N} \sum_{m,n=-l}^{l} \sum_{k=-l}^{l} c_{m,n}^l D_{m,k}^l(\mathbf{x}) \int_{G} \overline{D}_{n,k}^l(\mathbf{u}) \mathbf{f}(\mathbf{u}) d\mathbf{u}$.

So, $J_N(f) \in \Pi_N$.

Theorem 2.10. If $f\in L_p(SO(3)), 1\leqslant p\leqslant +\infty,$ then

$$\|\mathbf{f} - \mathbf{J}_{\mathbf{N}}(\mathbf{f})\|_{\mathbf{p}} \leq C\omega_{\mathbf{r}}(\mathbf{f}, \frac{1}{N})_{\mathbf{p}},$$

with C > 0 being a constant independent of f and N. *Proof.*

$$f(\mathbf{x}) - J_{N}(f, \mathbf{x}) = \int_{G} \phi_{N,s}(\mathbf{u}) \sum_{j=0}^{r} (-1)^{j-1} {r \choose j} f(\mathbf{u}^{-j}\mathbf{x}) d\mathbf{u}.$$

By Minkowski's inequality and (2.3),

$$\begin{split} \|f - J_{N}(f)\|_{p} &\leq \int_{G} \phi_{N,s}(\mathbf{u}) \|\Delta_{h}^{r} f\|_{p} \, d\mathbf{u} \\ &\leq \int_{G} \phi_{N,s}() \omega_{r}(f, |\mathbf{u}|)_{p} \, d\mathbf{u} \\ &\leq \omega_{r}(f, 1/N)_{\mathbf{u}p} \int_{G} \phi_{N,s}(\mathbf{u})(1 + N|\mathbf{u}|)^{r} \, d\mathbf{u} \\ &\leq C \omega_{r}(f, 1/N)_{p}. \end{split}$$

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